

# Realization of Tunable Channel Drop Filter Based on Square Photonic Crystal Ring Resonator

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#### Abstract

We have proposed a tunable two dimensional photonic crystal channel-drop filter (CDF) using ring resonators with suitable quality factor and transmission efficiency; we investigate parameters which a ffects resonant wavelength in this CDF, such as dielectric constant of inner, coupling, adjacent and whole rods of the structure and radius of inner rods. Dropping efficiency at the resonance and quality factor (Q) of our proposed CDF are 100% and 196, respectively. The footprint of the proposed structure is about  $125\mu m^2$ .

Keywords: Filter, Photonic crystal, Ring Resonator, Wavelength.

# **INTRODUCTION**

Photonic c rystals (PhCs), a lso k nown a s ph otonic bandgap (PBG) m aterials, c an c ontrol the s pontaneous e mission a nd the propagation of el ectromagnetic (EM) w aves [1-3]. Due t o existence of P BG, P hCs h ave ap plications in d ifferent ar eas o f optical en gineering s uch as o ptical f ilters [4], p olarizers [5], polarization s plitters [6], a nd demultiplexers [7] w hich m ay ultimately p ave the way for p hotonic in tegrated c ircuits (PICs). By engineering the photonic band gap, the confinement of light in given wavelengths can be tuned. It means that, we can control the transmittance and r eflectance w avelengths i nterval b y ad justing the bandgap of each PhC structure. In other words, PhC structures suggest h igh s pectral s electivity th at is necessary for f ilter designing [3]. F iltering i s the o peration e nabling us to e xtract from one w aveguide one f requency a nd to s end it t o a nother waveguide [8]. R ing r esonators have be en us ed as the building

blocks for the synthesis of high-order optical filters [9,10]. The distinctive feature of our proposed structure is that the resonant wavelength of ring resonator is tunable. Also design parameters for t his r ing r esonator c an be t he di electric c onstant of t he coupling, adjacent and whole rods of the structure. Compared to point-defect or line-defect PhC cavities, PhC ring resonators offer scalability in s ize a nd f lexibility in mode d esign d ue t o th eir multimode n ature [11]. W e u sed this element to ach ieve a n ew type o f C DF w ith h igh n ormalized t ransmission (100%) a nd acceptable qua lity factor (over 196) in 1500-1600*nm* window. The new ring resonator introduced in this study can be used as the basic element for other devices as well.

In t his pa per effects o f r ing p arameter o n t he r esonance wavelength a nd transmission s pectrum of t he s tructure a re investigated. The proposed device provides a possibility of optical channel dr op f ilter a nd c an b e us ed i n t he f uture p hotonic integrated circuits. Channel Drop Filter Design Using Photonic Crystal Ring Resonator



Figure 1. (a) A channel drop filter

A typical ring resonator obtained by removing a ring shape of co lumns f rom a s quare l attice o f d ielectric r ods in ai r background is displayed in Fig. 1(a). The dielectric rods have a dielectric constant ( $\varepsilon_r$ ) of 11.56, and radius r=0.2a is located in air, where a is a lattice constant. To m inimize the effect of counter propagating mode resulting from back-reflections at the sharp corners of the ring, we add one scatterer rod at each corner at half lattice constant as shown in Fig. 1(a). This additional rod at each corner act s as a r ight an gled r eflector r educing t he back-reflection a tt he c orresponding c orner [ 12,13]. For improving t ransmitted pow er t o p ort B, and ob taining more coupling e fficiency, r adius of c oupling, s catterer an d a djacent rods is set to 1.05r, 1.05r and 0.95r, respectively, while r is the radius of other rods [7,12]. By putting a waveguide be side the ring r esonator, t he w aveguide a t i ts r esonant f requency c an be coupled to the ring resonator to trap the electromagnetic energy propagating in the waveguide and localize it in the ring resonator. In ot her w ords, t he r ing r esonator dr ops l ight f rom t he t op waveguide and sends it to the bottom waveguide. This structure consists of an input waveguide and two output channels labeled as A and B.

In normal case when the ring does not resonate, whole of power remains in the horizontal bus waveguide and goes through port A. A t r esonance c ase, m ost of the power t ransfers t o the bottom waveguide and goes through port B.

# SIMULATIONS AND RESULTS



Figure 1. (b) Band diagram of our proposed CDF for  $\varepsilon$ =11.56.

In this structure, as shown in Fig. 1(b), band gap exists from  $a/\lambda=0.29$  to  $a/\lambda=0.42$  for TM polarization (in which the magnetic field is i n propagation p lane a nd th e e lectric f ield is perpendicular), w here  $\lambda$  is the wavelength i n f ree s pace. The spectrum o f th e p ower tr ansmission is o btained w ith f inite difference t ime d omain (FDTD) m ethod. FDTD m ethod is the most f amous method f or PhC analysis [14]. F DTD is a tim e domain s imulation m ethod f or solving M axwell's e quations i n arbitrary materials an d g eometrics. T he b asic principle is t o substitute t he c url e quations a nd partial t ime d ifferential w ith finite cen tral d ifferences i n s patial d omain and t ime d omain respectively.

Berenger's p erfectly matched l ayers (PML) ar el ocated around the whole structure as absorbing boundary condition [15]. The result of the FDTD simulation for this CDF that shows the normalized optical power transmissions of the structure is shown in Fig. 2. As shown in Fig. 2, the wavelength of  $\lambda$ =1543nm of the input port is removed from the upper waveguide and transmitted to t he p ort 'B'. T he tr ansmitted p ower e fficiency in th is wavelength i s a bout 1 00%. T he v alue of Q for t he pr oposed structure is obtained 196.

Figs. 3 (a) and (b) shows the electric field distributions of the structure proposed in Fig. 1(a) for two different wavelengths,  $\lambda_1=1543nm$ , t he r esonant a nd  $\lambda_2=1500nm$ , t he of f-resonance wavelength. At resonant wavelength, the field of the waveguide

fully c ouples t o t he r ing, w hereas a t of f-resonance i t does not couple and continues to propagate through the top waveguide.



Figure 2. Optical power transmission characteristic of our proposed CDF



Figure 3. (a) Electric field intensity of the ring resonator's non-resonant

wavelength at 1500nm and



(b) e lectric f ield i ntensity of t he r ing r esonator's r esonant wavelength at 1543*nm*.

The salient feature of this structure is that by varying the ring's parameters, the resonant wavelength can be tuned. In next section, the effect of varying some parameters on ring resonator performance will b e s tudied. S ection 3.1 d escribes effect of varying dielectric constant of rods, Section 3.2 describes effect of varying radius of rods and Section 3.3 describes effect of varying fill factor in resonance wavelength.

#### Varying Dielectric Constant of Rods

One of t he m ost i mportant f eatures of a ny f ilter is its tunability. Here we investigate parameters which a ffect resonant frequency in photonic crystal CDFs. F irst of all, we change the dielectric c onstant of t he w hole r ods. A s s een i n F ig. 4, t he proposed s tructure, w hen s imulated with the di fferent d ielectric constants equal to 10.24, 11.56 and 12.96 can select wavelengths of 1526*nm*, 1543*nm*, and 1561*nm*, respectively.



**Figure 4.** Power transmission spectra of our proposed CDF for different dielectric constants of whole rods.

As shown in Fig. 4, by raising the dielectric constant, the resonant w avelength of t he d evice is i ncreased accordingly. I n other words, a red shift occurs in resonant wavelength. We can create t he d ifferent d ielectric constants i n r eality b y u sing thermo-optic o r el ectro-optic m aterial. W e c an u set he thermo-optic effect caused by two-photon absorption (TPA) in *Si* to c ontrol t he r esonator's r effactive i ndex t hrough t he heat generated b y o ptically p roduced car riers; al so w e u tilize electro-optic m aterials which ch ange t heir r effactive in dexes in response to external electric field [16].

With localized change in inner rods' dielectric constant, the resonant wavelength can be tuned. This leads to a tunable CDF. Fig. 5 shows the normalized power transmissions of the structure with th ree d ifferent d ielectric constants o f i nner r ods,  $\varepsilon_r$ -0.4,  $\varepsilon_r$  and  $\varepsilon_r$ +0.4. A s s hown i n F ig. 5, our proposed s tructure, w hen simulated with th e d ifferent d ielectric c onstants o f in ner r ods equal to  $\varepsilon_r$ -0.4,  $\varepsilon_r$  and  $\varepsilon_r$ +0.4, can select wavelengths of 1540.8*nm*, 1543*nm*, and 1544.2*nm*, respectively.

In similar way, we can show if the dielectric constant of coupling or adjacent rods are changed, the resonant wavelength and so the whole CDF's characteristic spectrum will be changed similar to those of Fig. 5.



**Figure 5.** Power transmission spectra of our proposed CDF for different dielectric constants of inner rods.

#### Varying Radius of Rods

In this section we probe another parameter which affects on resonance behavior of the CDF. We study varying the radius of inner r ods f rom 0. 92r to 1. 1r. This l eads t o s hift of r esonant wavelength t o h igher v alues a s s hown i n F ig. 6. F or best resolution and tuning the filter, varying the dielectric constant of the inner rods as well as coupling rods can be used [16].

As shown in Fig. 6, it is obvious that when the radius of inner r ods of C DF i ncreases, the cu rves s hift t o r ight (greater wavelengths), which m eans t hat t his s tructure i s t unable f or desired w avelengths. A s s hown in F ig. 6, it is obvious that resonant wavelength i ncreases a s radius of inner rods i ncreases from 0.92r to 1.1r. As shown in Fig. 6, with different radius of

inner rods equal to 0.92r, r and 1.1r, we obtained wavelengths of 1538.6*nm*, 1543*nm* and 1550.5*nm* in port B, respectively.



**Figure 6.** Power transmission spectra of our proposed CDF for different radius of inner rods.

#### Varying Fill Factor of Structure

In this section, we investigate varying the fill factor of the structure f rom 0.  $98 \times r/a$  to 1.  $02 \times r/a$ . T his le ads to s hift o f resonant wavelength t o hi gher v alues a s s hown in F ig. 7. Our numerical s imulations r eveal t he d ependence of t he C DF's characteristic o n d ifferent r ing p arameters s uch as d ielectric constant of i nner, c oupling, a djacent a nd w hole r ods. A lso varying the fill factor as well as changing the dielectric constant of rods affects the CDF's characteristic.



Figure 7. Power transmission spectra of our proposed CDF for different fill factors.

# CONCLUSIONS

A t unable 2D phot onic c rystal C DF ba sed on r ing resonators had been introduced and investigated through FDTD method. By using a single ring resonator, we obtained the output power efficiency close to 1 00%. We investigated the effects of ring's pa rameters s uch a s i nner r ods' r adius; i nner, c oupling, adjacent and whole rods' dielectric constant and fill factor on the resonance w avelength. I t was s hown t hat t he r esonance wavelength of C DF has be en t uned by varying t his pa rameter appropriately. We have shown that there is flexibility in design of the CDF with photonic crystal ring resonators.

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